Estimating Balance Velocities Using GIS-based Techniques

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By

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ABSTRACT

We estimate Antarctic ice flow balance velocities, which are the average speeds that ice flows through ice cross sections if it is assumed that the input ice of these cross sections is equal to the output of these cross sections. The balance velocities are calculated using the OSU Digital Elevation Model (DEM), most recent ice accumulation rate and ice thickness data for Antarctic. We choose flux algorithms that allow efficient and accurate computation. The effective merging of digitized flow stripes from satellite images with our modeled flow directions mitigates the problem of problematic flow direction estimates in flat areas. In addition standard statistical theory of error propagation is used to estimate errors for balance velocities from input data measurement errors. Our new model compares favorably with previous research. Comparison between balance velocities and flow stripes suggests that fast glacier flow extends deeper into Antarctica than determined by flow stripe data above.

Key words: Hydrological modeling, Balance velocity, Antarctica

Dedicated to my parents

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CHAPTER 1

INTRODUCTION

Polar ice sheets spread and thin under their own weight in a fashion dictated by the constitutive relationship of ice and the external forces acting on their sides. Under the simplest of conditions, these processes would result in ice sheets of parabolic surface profiles and some axial symmetry (Paterson, 1994). Flow through such a simple ice sheet would be governed by the amount of snow accumulated on the surface and any mass lost from the base, terminus, or surface melting. In reality, the shape and flow of ice sheets are much richer features because of the interplay of complex boundary conditions and even, perhaps, spatial variability in the creep behavior of the ice through crystal reorientation. This interplay results in a variety of flow styles across both of the polar ice sheets. These may be characterized by the flow of the interior ice sheet, the surface of which most closely mimics the predicted near parabolic shape, the flow of ice streams, which are huge rivers of ice that meander through the interior ice sheet, and the marginal ice shelves that are vast, nearly flat, slabs of ice floating on the polar ocean.

Ice sheet shape, surface velocity and surface mass balance provide important clues about ice sheet dynamics and the ice sheet equilibrium state, so acquiring these

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observations has been the focus of over 50 years of research. Surface shape can be measured using a variety of altimetric techniques. The shape of the basal interface is commonly measured with ice sounding radar. Surface accumulation rate is traditionally measured using in situ methods but several investigations now suggest that remote sensing techniques may be applicable in some circumstances (Bolzart and Jezek, 2000). Until very recently, velocity has been measured using in situ data, but feature retracking on satellite images and radar interferometry are showing that this observation may best be done from space (Gray *et al.*, 2001).

Given these data, we now pose a question to provide some context for interpreting the observations. The simplest question is to ask whether or not the ice sheet is in steady state. That is, we ask whether or not the ice sheet is thickening or thinning. One way to answer that question is to consider what the physical properties of the ice sheet would have to be if it were in steady state. A traditional approach has been to take information on ice sheet surface and basal topography, surface and basal accumulation rate, and calculate the predicted surface velocity using mass continuity. The calculated velocities are known as balance velocities because they represent the ice sheet surface speeds that would result if the amount of snow added annually to the surface equaled the amount of mass lost from a particular area through advection or melting.

This thesis presents a new approach for computing Antarctic ice sheet balance velocities. It builds on previous research (Budd and Carter 1971; Budd and Allison 1975;

Smith and Budd, 1981; Budd *et al.* 1982; Radok *et al.* 1982; Budd and Warner, 1996; Bamber *et al.* 2000, 2001) by investigating several new algorithmic approaches and utilizing the best available data. It also attempts to make use of the recently completed Radarsat Antarctic Mapping Project SAR mosaic by incorporating derived vector data on flow direction into balance velocity estimation.

The results of this analysis are compared to measured surface velocities for the Lambert Drainage Basin and for several of the ice streams which drain into the Filchner Ice Shelf. The results suggest that only one of the Filchner ice streams may be out of balance (within the limitations of the data).

The balance velocity results will also be used to constrain velocities that will be measured using Radarsat interferometric data collected during the Modified Antarctic Mapping Mission of 2000. The model results will be used in areas of low surface speeds (less than 20 m per year) and where no other data are available. Because using the balance velocities biases the interferometric SAR (InSAR) results towards ice sheet which is in equilibrium, care will be taken to avoid propagating the model results into regions where non-steady state conditions (such as ice streams) may exist.

The concept of balance velocity is introduced in the second chapter and previous work on balance velocity and related work in hydrology are reviewed. Major problems associated with this work are summarized. In the third chapter, previous flow direction

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algorithms are discussed and a new technique to mitigate flow direction issues in low slope area is presented. In the fourth chapter, previous flux estimation algorithms are compared and a modified flux calculation approach is designed which is robust against DEM errors. In the fifth chapter, Antarctic DEM, accumulation rate and ice thickness data sets are utilized for the balance velocity calculation. Finally in the sixth chapter, conclusions are drawn.

CHAPTER2

BALANCE VELOCITY

2.1 Introduction

Hydrological models of water flow using DEMs are well known (Peucker and Douglas, 1975). Essentially, the models compute downslope directions and incorporate knowledge of water sources to estimate flow patterns and discharge. A similar concept, balance velocity, has been applied in glaciology (Budd *et al.* 1971). Balance velocities are the average speeds of ice flow through ice cross-sections if it is assumed that the ice input to these cross-sections is equal to the output (Figure 2.1). The following notation is used:

- *a* Accumulation rate with unit of kg/m^2a
- A Surface area with unit of m^2
- V_i Balance velocity of input ice with unit of m/a
- h_i Ice thickness of input ice flux gate with unit of m
- w_i Width of input ice flux gate with unit of m
- V_0 Balance velocity of output ice with unit of m/a

- h_0 Ice thickness at output flux gate with unit of m
- w_0 Width of output ice flux gate with unit of m
- ρ Constant ice density with unit of kg/m^3

Figure 2.1 shows that there are two sources of ice into the volume element which is under surface area A: snow accumulated on the surface and ice advected from higher elevations. The total volume flux of snow from accumulations (a A) equals the accumulation rate times the surface area. It is assumed that the ice density is constant (917 kg/m^3). The advected ice flux ($V_ih_iw_i\rho$) equals the balance velocity of the input to the volume element times the cross sectional area and the ice density. There is no input flux to the volume from the sides if the sides are taken parallel to flow lines. The only discharge of ice from the volume is through the downstream gate and this ($V_0h_0w_0\rho$) equals the cross sectional area times the balance velocity (to be solved for) and the ice density. An additional flux (either accumulation or discharge) occurs at the base of the volume element (melting or freezing), but this possibility is ignored in this analysis.

Since it is assumed that the ice sheet is in equilibrium, the total input ice is equal to the total discharge.

$$aA + V_i h_i w_i \rho = V_0 h_0 w_0 \rho \tag{2.1}$$

Then the balance velocity (V_0) is:

$$V_0 = \frac{a A + V_i h_i w_i \rho}{h_0 w_0 \rho}$$
(2.2)



Figure 2.1 Concept of balance velocity

2.2 Previous Methods to Estimate Balance Velocity

It is easy to understand the concept of balance velocity. However, many researchers continue to try to design algorithms for estimating balance velocity.

While balance velocity is conceptually easy to understand, algorithms for efficiently and accurately calculating it involve numerical and / or physical simplifications. Budd *et al.* (1971) proposed a flowline-type technique. In their model, they examined a block of ice bounded by digitized flowlines. Then, the net surface accumulation was calculated. This amount of ice was considered as the total input ice for a specific flux gate. Finally, the average balance velocity at this flux gate equaled the total input ice divided by the flux gate area. This technique was subsequently used by other researchers (Budd and Carter 1971; Budd and Allison 1975; Smith and Budd, 1981; Budd *et al.* 1982; Radok *et al.* 1982). Although this technique reflects the concept of balance velocity, its manual processing makes it time consuming to calculate. At places where flowlines are hard to determine, the result is less accurate than automatic numerical models designed by others. Budd and Smith (1985) devised an automatic gridded technique that has been widely used in a number of applications (Budd *et al.* 1986; Radok *et al.* 1986, 1987; Budd and Jenssen 1987, 1989; Mavrakis 1993). In the first step of this technique, summarized in Budd and Warner (1996), flow direction is calculated by estimating the direction of steepest slope using a digital elevation model (DEM). In Figure 2.2, cell $x_{i,j}$ has four cardinal nearest neighbors: $x_{i-1,j}$, $x_{i,j+1}$, $x_{i+1,j}$ and $x_{i,j-1}$. The surface elevations for them are denoted $E_{i,j}$, $E_{i-1,j}$, $E_{i,j+1}$, $E_{i+1,j}$ and $E_{i,j-1}$. With this, the horizontal slope component (α_x) and the vertical slope component (α_y) for cell $x_{i,j}$ are:

$$\alpha_x = (E_{i-1,j} - E_{i+1,j}) / 2\Delta x \tag{2.3}$$

$$\alpha_{y} = (E_{i,j-1} - E_{i,j+1}) / 2\Delta y \tag{2.4}$$

The flow direction (θ) is given by

$\sin\theta =$	$\frac{\alpha_y}{\alpha},\cos\theta$	$\theta = \frac{\alpha_x}{\alpha}$		(2.5)
		<i>x</i> _{<i>i</i>-1,<i>j</i>}		
	$x_{i,j-1}$	$X_{i,j}$	$x_{i,j+1}$	
		$x_{i+1,j}$		

Figure 2.2 Grid cell and its neighbors

Next, fluxes are identified for a particular cell. The total ice discharges from cell $x_{i,j}$ are denoted $\psi_{i,j}^{(out)}$; the inputs of ice to cell $x_{i,j}$ from neighboring cells are denoted

 $\psi_{i,j}^{(in)}$. $A_{i,j}$ is the accumulation rate at cell $x_{i,j}$; Δx is the width of the cell and Δy is the height of the cell. The discharge from the cell is then given by:



 $\boldsymbol{\psi}_{i,j}^{(out)} = A_{i,j} \Delta x \Delta y + \boldsymbol{\psi}_{i,j}^{(in)}$ (2.6)

Figure 2.3 Flux partition

 $\psi_{i,j}^{(in)}$ is calculated from discharge from the neighboring cells of $x_{i,j}$. A simple partitioning scheme partitions the flux by dividing the flux magnitude by the normalized form of the directional cosines. For example, in Figure 2.3 $\psi_{i,j}^{(in)}$ comes partly from two neighbor cells: the y component of $x_{i,j+1}$ ($\psi_{i,j+1}^{(out)Y}$) and the x component of $x_{i+1,j}$ ($\psi_{i+1,j}^{(out)X}$), defined as follows.

$$\psi_{i,j+1}^{(out)Y} = \frac{\left|\sin \theta_{i,j+1} \middle| \psi_{i,j+1}^{(out)} - \frac{1}{\left|\sin \theta_{i,j+1} \right| + \left|\cos \theta_{i,j+1} \right|}$$
(2.7)

$$\psi_{i+1,j}^{(out)X} = \frac{\left|\cos\theta_{i+1,j}\right| \psi_{i+1,j}^{(out)}}{\left|\sin\theta_{i+1,j}\right| + \left|\cos\theta_{i+1,j}\right|}$$
(2.8)

The $\sin \theta_{i,j+1}$, $\cos \theta_{i,j+1}$, $\sin \theta_{i+1,j}$ and $\cos \theta_{i+1,j}$ terms in equation (2.7) and (2.8) are calculated from equation (2.5).

The calculation is implemented by sorting the cells by descending elevation. From the highest cell to the lowest cell, the output flux for each cell is calculated so that the total output flux of every upper input cell is always known from previous calculations.

Implementing this gridded technique, Budd and Warner (1996) created an ice flux distribution map for Antarctica. Based on this technique, Bamber *et al.* (2000, 2001) calculated balance velocity maps for Greenland and Antarctica.

2.3 Review of Related Methods in Hydrology

Previous research on hydrological models provides useful ideas for improving on balance velocity models. Here several approaches are discussed which have a bearing on the work conducted as part of this investigation.

Lea (1992) designed a best fit plane algorithm to get flow direction from a DEM using the principle of least squares. He fitted a plane to four corners of a cell. The corner heights are calculated from the average of spot heights of four neighbor cells to each corner (Figure 2.4). For example,

$$C_{ne} = \frac{E_{i-1,j-1} + E_{i-1,j} + E_{i,j-1} + E_{i,j}}{4}$$
(2.9)

"ne", "nw", "se", "sw" represent four corners of center cell $x_{i,j}$. C_{ne} , C_{nw} , C_{se} and C_{sw} are the heights at the corners.

	<i>x</i> _{<i>i</i>-1,<i>j</i>-1}	$X_{i-1,j}$	$x_{i-1, j+1}$	
	nv nv	V 1	ne	
	$x_{i,j-1}$	$X_{i,j}$	$x_{i,j+1}$	
	SW	/ .	se	
	$X_{i+1, j-1}$	$X_{i+1,j}$	$x_{i+1,j+1}$	
Figure 2.4 Fitting plane algorithm				

Costa-Cabral and Burges (1994) devised a two-dimensional flux partition scheme. They also proposed that partitioning should be proportional to areas defined by flow directions rather than proportional to flow directions.



Figure 2.5 Two-dimensional partition scheme

Figure 2.5 illustrates partitioning proportional to areas. For cell $x_{i,j}$, the black arrow represents its flow direction (θ). The dashed line (CE) is the flow line passing through the lowest corner point (E) and establishes the boundary for mass partitioning. The blue arrows represent the mass in polygon (ABCE) which flows into cell $x_{i,j-1}$ while the red arrow represents the mass in triangle (CDE) which flows into cell $x_{i+1,j}$. The fluxes partitioned to cell $x_{i,j-1}$ and to cell $x_{i+1,j}$ according to the scheme of Costa-Cabral and Burges (1994) are:

Area of ABCE:
$$1 - \frac{|tg(\theta - 90)|}{2} = 1 - \frac{|ctg\theta|}{2}$$
 unit (2.10)
Area of DCE:
$$\frac{|tg(\theta - 90)|}{2} = \frac{|ctg\theta|}{2}$$
 unit (2.11)

Most partitioning schemes lead to a false diffusion of mass. This comes in part from the partitioning scheme and in part because of the use of Cartesian grids. Costa-Cabral and Burges (1994) designed a scheme to avoid this problem. The scheme relies on detailed flow directions to define the contribution area bounded by flow lines (hence stopping diffusion across flow boundaries). The scheme is complex and Burges (2001) has pointed out that it is computationally difficult to implement. Moreover, small errors in flow directions will create significant errors in contribution area, making the model unstable when implemented over large areas. The model is valuable for simple terrains and for better understanding the effects of diffusion on conceptually less sophisticated partitioning algorithms.

2.4 Limitations of previous approaches

Balance velocity calculations begin by estimating the flow direction, taken to be along the steepest slope. Then, accumulation rate and flow direction are combined in the model to calculate discharge flux. Dividing the flux by flux-gate area, the balance velocity is finally specified. All flow direction algorithms require DEM information about neighboring cells. If the height difference of neighbor cells is small or zero (to the precision of the original methods and gridding schemes), the algorithms cannot accurately determine the flow direction for the cell. This problem is particularly acute on polar ice shelves, which are large floating slabs of ice with typical surface slopes less than 0.01°. Another problem peculiar to glacier flow is that the ice is assumed to flow down the mean surface gradient. Superimposed on the mean slope are local variations arising from longitudinal stresses. Most models deal with longitudinal variations by averaging the surface elevation over a distance equal to about 10-20 times the ice thickness (Paterson, 1994; Budd and Warner, 1996; Bamber *et al.* 2001).

Diffusion is a rarely discussed problem in balance velocity calculations. Essentially the partitioning schemes allow ice to flow across flow lines. For example, suppose a planar topography tilting from Northwest to Southeast. Figure 2.6 shows how the flux at cell A flows to its downstream cells. The region between two dashed lines is bounded by flow lines. Simple partitioning schemes which allocate ice per the Cartesian grid diffuse ice outward from this band. In Figure 2.7, the small shaded squares represent the proportion of the flux from cell A passing into the cells. The diffusion effect is quite evident.



Figure 2.6 Problem of flux model (Costa-Cabral and Burges, 1994)



Figure 2.7 Problem of flux model (Costa-Cabral and Burges, 1994)

Since low slope areas are typical for Antarctica, the problem that flow directions are affected by DEM errors is important for the Antarctic balance velocity model. Diffusion distorts flux estimation. Our algorithm, discussed in the following chapters, attempts where possible to correct some of these limitations. Where not possible, we at least attempt to quantify the error introduced by the limitations. In Chapter 3, a solution is proposed which mitigates some of the errors associated with low slope terrain. In Chapter 4, the degree of diffusion is quantified and discussed.

CHAPTER 3

FLOW DIRECTION

3.1 Introduction and previous work

Flow directions are used to map drainage networks as a first step in calculating mass discharges. Digital elevation models (DEMs) are the basis for deriving flow directions. There are numerous schemes for estimating flow direction from a DEM. A simple scheme (D8), assigns eight possible flow directions to each cell (O'Callagham and Mark, 1984). All of the discharge is assigned to only one flow direction of the eight possible directions. A refinement is the multiple flow direction model (Quinn et al. 1991; Freeman, 1991) in which all eight downslope cells are assigned as possible flow directions (as might be envisioned at the termination of a ridge). Finally fitting plane algorithms (Lea, 1992; Costa-Cabral and Burges, 1994; Tarboton, 1997) estimate the direction of maximum slope and assign that direction to be the sole flow direction.

In Chapter 2, Lea's (1992) fitting plane algorithm was introduced. Costa-Cabral and Burges (1994) used four diagonal cell heights to fit a plane with a modified version of Lea's method and assigned the steepest slope of the plane as the flow direction (Figure 3.1). The x and y components of the slope are:

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$$\alpha_x = \frac{E_{i-1,j+1} - E_{i-1,j-1} + E_{i+1,j+1} - E_{i+1,j-1}}{2}$$
(3.1)

$$\alpha_{y} = \frac{E_{i-1,j+1} - E_{i+1,j+1} + E_{i-1,j-1} - E_{i+1,j-1}}{2}$$
(3.2)

i-1,j-1		i-1,j+1
	i,j	
i+1,j-1		i+1,j+1

Figure 3.1 Fitting plane algorithm of Costa-Cabral and Burges

Because only three points are needed to describe a plane, Tarboton (1997) proposed an algorithm to fit a plane to three cell heights. In Figure 3.2, letter A – H and O are the center points of the grid cells. The facet OAB is the plane fit to cell heights of $x_{i,j}$, $x_{i,j+1}$ and $x_{i-1,j}$. There are eight fitting facets for cell $x_{i,j}$. For facet OAB, the x and y components of the facing direction of the facet are:

$$\alpha_{x} = \frac{E_{i,j} - E_{i,j+1}}{cellsize}$$
(3.3)
$$\alpha_{y} = \frac{E_{i,j+1} - E_{i-1,j+1}}{cellsize}$$
(3.4)

Tarboton calculated the facing directions for the eight facets of the cell $x_{i,j}$. Comparing the steepest slope magnitudes of these eight facets, the facing direction of the biggest slope facet is assigned as the flow direction for the cell $x_{i,j}$.



Figure 3.2 Fitting plane algorithm of Tarboton

The flow direction algorithm used by Budd and others (1996) is similar to the Costa-Cabral and Burges method but uses the cardinal rather than diagonal points about the test cell. Referring to Figure 3.2, it fits a plane to the four cells $x_{i-1,j}$, $x_{i,j+1}$, $x_{i+1,j}$ and $x_{i,j-1}$.

3.2 Fitting plane algorithm

For simple terrain, the above fitting plane algorithms calculate similar flow directions, but for complicated terrain or DEMs with noisy data, they create different results. For example, Figure 3.3 is a hypothetical DEM (the numbers in the grid cells are heights). The flow direction using Lea's (1992) algorithm is 123.7°, using Costa-Cabral and Burges (1994) it is 199.2°, using Budd (1996) it is 110.2°, using Tarboton (1997) it is 191.3°. Essentially the outlier value in the lower center cell biases results which are

derived from averaging multiple points. The outlier is either ignored or over-emphasized in schemes which only use diagonal or cardinal points.

7 13 14	9	11	12
	7	13	14
5.8 30 12	5.8	30	12

Figure 3.3 Hypothetical DEM

When we applied these four algorithms to the Antarctic DEM and compared the results with flow stripe directions which are the orientation of flow stripes digitized from image data (detail in next section), we found that the four algorithms behave similarly in terms of means and deviations of difference from flow stripe directions.

We did not find any conclusive quantitative criteria for selecting one approach over the other. Ultimately we used the Costa-Cabral and Burges algorithm because the ice sheet generally exhibits smooth variation, outliers were as best as possible removed in the original DEM generation, and the fit to the four most widely separated cells (diagonal cells about the evaluation cell) effectively resulted in additional spatial averaging of the surface slopes.

3.3 Flow stripes digitized from RAMP Radarsat Image

The flow-direction distribution map of Antarctica is created by applying the fitting plane algorithm to the Antarctic DEM. For reasons mentioned above, there are some flow regimes where the direction is unreliable. Ice streams are relatively narrow features (order of 50 km) that are difficult to capture with our grid cell spacing (20 km as discussed later). Ice shelves are very flat and slight errors in the DEM lead to large errors in calculated flow direction.

The problem of modeling flow in regions of low slope is discussed by Garbrecht and Martz (1997) who presented an approach that modifies flat surfaces to force the water to flow from higher to lower terrain. The approach assigns flow directions in flat areas that force convergence. This approach does not apply to our study because the flow through ice streams and on ice shelves may be either convergent or divergent. Liang and Machay (1997, 2000) designed an automated method to assign flow directions for flat areas using a T-test (difference of means).

We present an accurate and straight forward method here to mitigate the ice shelf and ice stream problems. Flow stripes in the fast ice streams and ice shelves are evident on the RAMP Radarsat imagery (1997). In regions where the ice has been in equilibrium, these are known to accurately reflect the flow direction. As part of a separate project, the stripes were manually digitized and converted into Arc/Info line coverage. Figure 3.4(a) shows the digitized flow stripe map from the Radarsat Image Figure 3.4(b). Figure 3.4(c) llustrates the accuracy of digitization in east Antarctica.

3.4 Flow directions from flow stripes

Our scheme to estimate flow direction from the digitized flow stripes (Sohn and Jezek, 1997) is as follows.

There are five steps in the process (Figure 3.5). In the first step, the line coverage of flow stripes is converted to x and y coordinates series. For example (Figure 3.6), line l_1 (ABCD) is converted to coordinates series from the starting point to the ending point: 1,4; 3,1; 5,0.5; 7,0.



Figure 3.4(a) Digitized flow stripes



Figure 3.4(b) Radarsat Image of Antarctica



Figure 3.4 (c) The accuracy of digitization



Figure 3.5 Process of flow stripes direction



Figure 3.6 Line point coordinate series

Next, the line coverage is separated into line segment coverage. For example, line l_1 is converted into segments AB, BC and CD. Following segmentation from the coordinate series, the orientation of each line segment of a flow stripe is calculated. For example the orientation angle for segment AB of line l_1 in Figure 3.6 is 123.7°. In the third step, a new attribute consisting of the facing angle for every segment is added into the line segment coverage database using Arc/Info. In the fourth step, the line segment coverage is gridded into a 20km cell-size grid. In the grid, the values of the cells are the facing angles of the line segments that are nearest to the centers of the cells. In the final step, the grid value is oriented to reference to the same digitization order because the different digitization orders of a line, such as digitizing from a starting point to an end or from the end to the starting point), give different facing angles of the same line segments. Figure 3.7 is the gray grid of flow stripe directions with blue flow stripes overlaid on it.


Figure 3.7 Grid of flow strip direction

3.5 Final flow direction map

We merge fitting-plane modeled flow-directions from the DEM with the flow stripe directions by preferentially selecting the flow stripe orientation (Figure 3.8). Figure 3.9 shows the difference between flow stripe directions and flow directions from the DEM using a fitting plane algorithm.

We do not calculate flow direction over the flat ice shelves (indicated as white in Figure 3.8). Very dense flow stripes can be drawn on ice shelves using an enhanced



Figure 3.8 Final flow direction grid with flow stripes



Figure 3.9 Direction difference between flow stripe directions and directions from DEM

composite Advanced Very High Resolutions radiometer (AVHRR) image (Fahnestock *et al.* 2000). In some cases, using the flow stripe directions on ice shelves can improve the flow direction map (Figure 3.8). However, it has been shown that flow stripes on ice shelves may be strongly influenced by past changes in ice shelf dynamics. Hence we choose not to include that information in our analysis.

CHAPTER 4

FLUX ESTIMATION

4.1 Introduction

In terms of glacial discharge, flux is the total amount of ice mass passing per unit time through a cross-sectional area whose surface is perpendicular to the flow direction. Essentially this represents continuity of mass. If the mass balance of part of the ice sheet is in equilibrium, then the flux entering a unit volume must equal the flux leaving that volume. Fluxes are calculated from knowledge about surface accumulation rate and ice thickness. Flow speeds and directions can be measured, or as done here, balance velocities can be computed. The magnitude of the balance velocity vector is determined by mass continuity and the direction of the balance velocity vector is derived from knowledge about surface elevation as described in the previous chapter.

Ideally, flux is treated as a continuous variable with the property that there is no flux across stream lines. Practical computation schemes require that the flux calculation be carried out on a discrete grid that then forces compromises on the analysis. The compromises arise when the discharge flux from a particular cell is partitioned into adjacent cells and when the summed discharges are converted into a flux. In this chapter,

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different partitioning schemes are presented. These are then incorporated into different flux algorithms for subsequent comparison.

4.2 Partition schemes

Given estimates of flow direction and mass discharge per cell, partition schemes are used to model how much mass flows from one cell into its neighboring cells. Budd and Warner (1997) partition the mass by dividing the flux magnitude by the values of normalized components in directions taken to be aligned with the grid direction. Other partition schemes distribute the mass according to the magnitudes of slopes or functions of slopes in the direction of neighboring cells (Quinn *et al.*, 1991; Freeman, 1991). Costa-Cabral and Burges (1994) proposed a scheme wherein partitioning is proportional to areas defined by segmenting the cell into two regions separated by the flow direction vector.

Figure 4.1 is an example illustrating differences between the partitioning schemes proposed by Budd and Warner (1996) and Costa-Cabral and Burges (1994). θ° is the flow direction. Suppose the flux at cell $x_{i,j}$ is 1 unit. The fluxes partitioned to cell $x_{i,j-1}$ and to cell $x_{i+1,j}$ according to the scheme of Budd and Warner (1996) are:

$$\frac{|\cos(\theta - 90)|}{|\sin(\theta - 90)| + |\cos(\theta - 90)|} = \frac{|\sin\theta|}{|\cos\theta| + |\sin\theta|} \text{ unit}$$
$$\frac{|\sin(\theta - 90)|}{|\sin(\theta - 90)| + |\cos(\theta - 90)|} = \frac{|\cos\theta|}{|\cos\theta| + |\sin\theta|} \text{ unit}$$

The fluxes partitioned to cell $x_{i,j-1}$ and to cell $x_{i+1,j}$ according to the scheme of Costa-Cabral and Burges (1994) are represented by equations 2.10 and 2.11.



Figure 4.1 Partition schemes (The upper right cell illustrates the Costa-Cabral and Burges method. The lower left corner illustrates the Budd and Warner method.)

4.3 Flux algorithms

Whereas partition schemes model how one cell discharges mass into neighboring cells, flux algorithms model how all the cell contributions form the drainage system. Again, different authors have proposed a variety of schemes. For example, Budd and Warner begin by sorting the DEM from the highest DEM cell to the lowest one to cover the entire terrain. There are some algorithms that do not require special processing order, such as the DEMON-downslope algorithm (Costa-Cabral and Burges, 1994), DEMON-upslope algorithm (Costa-Cabral and Burges, 1994) and upslope area calculation (Mark, 1988). These algorithms are less sensitive to DEM errors by not using DEM order in processing. Because we introduced the Budd and Warner algorithm in Chapter 2, DEMON algorithms and others are explained in the following sections.

4.3.1 Upslope area calculation (Mark, 1988)

This algorithm is a recursive procedure. First it identifies cells discharging into the reference cell. If the discharge fluxes of the cells are known, then the flux of the reference cell is calculated. If the flux from a potential discharge cell is unknown because the computation has not proceeded far enough along, the calculation proceeds to find the discharge cells for that new reference cell. The calculation proceeds iteratively until all of the discharges feeding the original reference cell are determined.

We applied Mark's algorithm in the flux calculation using Antarctica DEM and ice accumulation rate and compared the result with Budd and Warner algorithm (1996). We found that the algorithm overestimates the contributing area by about 20% in many places, though the shape of the ice stream networks are reasonable. Mark's algorithm uses flow directions to trace the contributing areas. The 20km cell-size flow direction grid cannot define the flow lines very accurately, so small errors in the flow directions create erroneously bigger contributing areas.

4.3.2 DEMON-Upslope algorithm (Castal-Cabral and Burges, 1994)

This algorithm calculates the contributing areas for every cell by tracing the flow lines passing through cells which define the boundary of contributing areas. The tracing stops until the ridges or the terrain boundary are encountered. Then the sums of the mass accumulated over the contributing area correspond by definition to the fluxes through every cell. Figure 4.2 shows the contributing areas for cell A with gray shading. The sum of the mass accumulated over this region is the flux of cell A. There are two major

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problems with this algorithm. Because it relies on flow directions to define contributing areas, flow direction errors create large errors in contributing areas. When two flow lines do not terminate within one cell, it is hard to estimate contributing areas. Thus the algorithm is difficult to implement in a computer program.



Figure 4.2 DEMON-Upslope algorithm

4.3.3 DEMON-Downslope algorithm (Costa-Cabral and Burges, 1994)

In this algorithm, each cell is treated as an initial condition wherein the flux entering the cell is simply the surface accumulation rate (advected fluxes are not included in the flux of the cell chosen for the start of the calculation). The discharge flux from the initial cell is allowed to flow downslope as an advective term until it encounters cells that are sinks or leaves the terrain. Once the discharge from every cell is calculated into a so call influence matrix, the total mass passing through each cell is summed (both accumulation and advection terms). In this way, the total mass passing through a particular cell equals the total mass drained by every upstream cell.

The influence matrix idea is illustrated Figure 4.3(a)). The gray region is the influence matrix of the cell A. Every cell in the gray region receives advected mass from cell A in an amount determined by the partitioning scheme. The boundaries of the influence matrix are flow lines so that there is no diffusion of mass across the boundary. Figure 4.3(b) is an example of how the value for cell B in the influence matrix is determined by flow lines. The darker region in cell A is the mass that passes through cell B. Though the influence matrix determined by flow lines is conceptually very accurate, it is hard to write an efficient and accurate computer program (Burges, 2001).



Figure 4.3 (a) Influence matrix (Costa-Cabral and Burges, 1994)



Figure 4.3 (b) Calculation of influence matrix (Costa-Cabral and Burges, 1994)

4.3.4 Modified Costa-Cabral and Burges Algorithm

We modify the influence matrix calculation by adopting a simpler partition scheme. This allows for a tractable code while retaining the advantage that there is no special ordering of the cells in the calculation. The new algorithm does not build flow lines. It successively partitions mass from one cell into its downslope cells according to the partitioning scheme proposed by Costa-Cabral and Burges. Partitioning proceeds until the edge of the terrain is reached (Figure 4.3(c)). In the Figure, A is partitioned into B and C. Then B is partitioned into D and E; C is partitioned into E and F. The cells with the same colors are processed at the same time – we call the ensemble of cells the front line. Front lines are not determined manually. They arise automatically and somewhat artificially from the process order. As discussed later, the front lines contribute to the cardinal partitioning scheme problem discussed later in section 3 in detail.



Figure 4.3 (c) New influence matrix

4.3.5 Comparison of flux algorithms

We compare the modified influence matrix algorithm with the Budd and Warner (1996) algorithm.

Budd and Warner's algorithm depends on the DEM order, which means that the order of calculation is sensitive to DEM error. The influence matrix algorithm does not use the DEM order to process cells. So the influence matrix algorithm is initially less affected by the DEM errors in flat or near flat areas than Budd and Warner's algorithm

The influence matrix approach also yields a picture of the up and downslope areas contributing to a cell. This aids in visualizing the flux calculation and help the validation of flux calculation.

Both algorithms suffer from the diffusion problem. They are also sensitive to the cell size chosen and the manner in which flux is partitioned between cells. In the following section these problems are discussed in detail.

4.4 Limitations of the flux algorithm

We discuss three limitations of the flux algorithms associated with diffusion, cell size and partitioning schemes.

4.4.1 Diffusion problem

By diffusion we mean the artificial distribution of mass away from boundaries that would properly be identified as flow lines. We use a simple hypothetical terrain to examine the degrees of diffusion for different flux algorithms. Figure 4.4 (a) is the hypothetical DEM which represents a plane tilting from Northeast to Southwest. The flow direction from the DEM is 120° from North direction represented by the arrow. The accumulated mass is assumed to be uniform (10 units). Because DEMON's Downslope area algorithm (Costa-Cabral and Burges, 1994) tries to incorporate a flow band constraint, we use it as a control. The flux modeled by DEMON's Downslope area algorithm is shown in Figure 4.4 (b). The flux modeled by our influence matrix algorithm is shown in Figure 4.4 (c). Figure 4.4 (d) is the percentage difference between the influence matrix algorithm and the DEMON's Downslope area algorithm. In Northwest sector the differences are on average about 40% disregarding the boundary results, which means the influence matrix algorithm overestimates flux. The error is biased to

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Northwest direction because the flow direction 120° causes more flux to be partitioned in the west direction and less flux in the south direction. Budd and Warner's (1996) algorithm creates similar results as influence matrix algorithm for the simple terrain.



Figure 4.4 (a) Hypothetical DEM



Figure 4.4 (b) Flux distribution of DEMON's Downslope area algorithm

(Costa-Cabral and Burges, 1994)



Figure 4.4 (c) Flux distribution of influence matrix algorithm



Figure 4.4 (d) Comparison of flux distributions of influence matrix and DEMON's

Downslope area algorithms (unit: percentage of difference)

4.3.2 Cardinal partitioning scheme problem

The partitioning scheme relegates the flux from a cell into its cardinal downstream cells. For example, figure 4.5(a) is a valley terrain where cell $x_{i,j}$ flows to cell $x_{i-1,j+1}$. The numbers in cells are their heights. According to the cardinal point partitioning schemes, flux from cell $x_{i,j}$ is partitioned into cell $x_{i-1,j}$ and cell $x_{i,j+1}$, but these are higher than cell $x_{i,j}$. This results in a conflict which causes cell $x_{i-1,j+1}$ to lose mass from cell $x_{i,j}$.



Figure 4.5(a) Conflicts between cardinal partition scheme and process order

This kind of conflict effects Budd and Warner's algorithm when the incorporated flow stripe derived directions conflict with the maximum slopes derived from the original DEM. The influence matrix algorithm also has this problem when flow directions are curved. In Chapter 5, this problem is discussed further.

4.3.3 Cell-size problem

The DEM assumes the terrain within a cell is uniform. Smaller scale terrain than the cell-size cannot be modeled by DEM. Although the DEM has a minimum 200 m cell size, it is well known that ice flows in the direction of maximum surface slope averaged over 10–20 times the ice thickness. This results in a fairly coarse grid. For example, when we apply the flux algorithm to Antarctica DEM with 20km cell-size, some small streams such as upper Ice Stream B cannot be detected.

CHAPTER 5

ANTARCTICA ICE SHEET BALANCE VELOCITY

5.1 Introduction

In this chapter, we present a new our balance velocity model. We describe the data and any preprocessing needs. We go on to validate the model by comparing it with independent estimates of surface velocity. Finally, we compare estimated balance velocity with Radarsat flow stripes to answer questions about the dynamics of the ice sheet.

5.2 Data Description

We used four primary data sets in this analysis: the OSU digital elevation model of Antarctica (Liu *et al.* 1999); the BEDMAP (2000) basal surface topography model; surface accumulation rates (Vaughan *et al.* 1999); and RAMP flow stripes. We describe these in the following sections.

5.2.1 OSU DEM of Antarctica

We use the most recent elevation data contained in the DEM from Liu *et al.* (1999) (Figure 5.1). The surface digital elevation model (DEM) is complied using several topographic data sets which have rather disparate sampling intervals. Consequently, the interpolated product was uniformly resampled to 200, 400 and 1000 m post spacings – recognizing that in some areas this represents an over sampling of the available data. Similarly the quality and accuracy of the data varies due to data collection methods.



Figure 5.1 Hillshade of Antarctica DEM

5.2.2 Accumulation rate

Accumulation rate data for Antarctica is summarized by Vaughan *et al.* (1999). The product is available as an Arc/Info grid with 10km cell-size. The data are presented in units of $kg/m^2 y$ (Figure 5.2). It is assembled from over 1800 published or unpublished in situ measurements. The uncertainty of the data is approximately 5% (Vaughan *et al.* 1999).



Figure 5.2 Accumulation rate of Antarctica (Vaughan et al. 1999)

5.2.3 Ice Thickness

Ice thickness was compiled as part of the BEDMAP project (Lythe, *et al.* 2000). The data is provided in an Arc/Info grid with 5km cell-size. It is based on about 2 million ice thickness observations by 12 countries over the last five decades (Figure 5.3). The accuracy of ice thickness is different in different regions and ranges from 10 to 180 meters (Lythe, *et al.* 2000). A few areas, such as the Amery Ice Shelf, Ronne-Filchner Ice Shelf and Siple Coast, have good accuracy and excellent coverage (Lythe, *et al.* 2000). Other areas, such as large parts of East Antarctica, are covered by 50km spaced flightlines or with little to no data at all (there, ice thickness is based on gross interpolation or model results) (Lythe, *et al.* 2000). We estimated ice thickness errors according to known properties of the data collection methods (Figure 5.4).



Figure 5.3 Distribution of ice thickness data in BEDMAP compilation (BEDMAP, 2000)



Figure 5.4 Ice thickness errors

5.3 Balance Velocity Estimation

This section summarizes the balance velocity estimation, including DEM preparation, flow direction estimation, flux estimation and balance velocity estimation.

5.3.1 DEM preparation

We use the OSU DEM with the 1km cell-size in Arc/Info grid format. Ice flow directions are determined by terrain slopes averaged over distances of 10-20 times ice thickness (Paterson, 1994). This is supported by Bamber *et al.* (2001) who adopted an averaging scale of 20 times the ice thickness. Thus the DEM is filtered by using a running, Gaussian weighting window of the following form:

$$G_{i,j} = \frac{e^{\frac{-d_{i,j}^2}{4r^2}} - e^{-\frac{1}{2}}}{\sum_i \sum_j \left(e^{\frac{-d_{i,j}^2}{4r^2}} - e^{-\frac{1}{2}}\right)}$$
(5.1)

where $G_{i,j}$ is the weight of the gaussian window, $d_{i,j}$ is the distance between the evaluation cell $x_{i,j}$ and any other cell within the window range, and r is the radius of the window taken to be half of 20 times the ice thickness. At the margin of Antarctica, the DEM is smoothed by local mean to avoid edge effects.

If there are any residual sinks (that is a local depression from the mean slope), they are filled using the Arc/Info grid function. The function raises the elevation of the sink to match the lowest height of the eight neighboring cells. The filtered DEM is resampled into a 20km cell-size grid, which corresponds to the cell size used by Budd and Warner (1996).

5.3.2 Flow direction estimation

As explained in Chapter 3, flow direction is first derived from the DEM using the Costa-Cabral and Burges fitting plane algorithm. Then, we merge the flow direction grid with the flow stripe directions by preferentially selecting the flow stripe orientation (Figure 3.8). Incorporation of vector data mitigates the sensitivity of flow direction to DEM errors in low slope areas and to some extent, regions of converging flow.

5.3.3 Flux estimation

As discussed in Chapter 4, a modified version of DEMON-Downslope algorithm provides an efficient and reasonably accurate computational approach. The balance velocity model is based on this algorithm.

5.3.4 Balance velocity estimation

We calculate the flux for every grid cell. The flux gate area for every grid cell is the cell-size dimension multiplied by the ice thickness. We treat the flux as a scalar with a fractional part of the flux going in the x and y directions. The portion of the flux in the x direction (F_x) is related to an equivalent balance speed (V_x) as:

$$F_x = V_x \cdot H \cdot W \tag{5.2}$$

Where H is the ice thickness, W is the grid cell-size dimension.

The portion of the flux in the y direction (F_y) is related to equivalent balance speed (V_y) as:

$$F_{y} = V_{y} \cdot H \cdot W \tag{5.3}$$

and

$$F = F_x + F_y \tag{5.4}$$

In turn, we define the components of the velocity vector (V) as:

$$F_{x} = V_{x} \cdot H \cdot W = V \cdot \cos \theta \cdot H \cdot W$$
(5.5)

$$F_{y} = V_{y} \cdot H \cdot W = V \cdot \sin \theta \cdot H \cdot W$$
(5.6)

Still treating everything as a scalar we add these equations to get:

$$V = \frac{F}{H \cdot W \cdot (\sin \theta + \cos \theta)}$$
(5.7)

where F is the flux for a grid cell with unit of $m^3 y^{-1}$, θ is flow direction previously determined by steepest slope from DEM.

Because balance velocity (V) is the average velocity at the flux gate, the ice surface balance velocity (V_s) is set equal to the balance velocity divided by 0.89 (Budd and Warner 1996). The ratio takes into account the fact that the velocity decreases from the surface towards the bed.

Figure 5.5 is the final surface balance velocity map.



Figure 5.5 Final surface balance velocity map

5.4 Error estimation

The uncertainty in the balance velocity comes from measurement errors associated with the input data and from algorithm.

5.4.1 Random error from input data

The random errors of DEM, accumulation rate and ice thickness are estimated from the data source documentation. These errors propagate when they are used to calculate flow direction, flux and final balance velocity. We apply the standard statistical theory of error propagation (Parrat, 1961 and Taylor, 1982) to estimate the random error from input data. The theory assumes no correlation between variables. The flow direction is determined by (see Chapter 3):

$$\alpha_{x} = \frac{E_{i-1,j+1} - E_{i-1,j-1} + E_{i+1,j+1} - E_{i+1,j-1}}{2}$$
(5.8)

$$\alpha_{y} = \frac{E_{i-1,j+1} - E_{i+1,j+1} + E_{i-1,j-1} - E_{i+1,j-1}}{2}$$
(5.9)

$$\theta = \tan^{-1}(\frac{\alpha_y}{\alpha_x}) \tag{5.10}$$

So the standard deviation of flow direction $(d\theta)$ is:

$$d^{2}\theta = \frac{d^{2}z}{\left(E_{i-1,j-1} - E_{i+1,j+1}\right)^{2} + \left(E_{i-1,j+1} - E_{i+1,j-1}\right)^{2}}$$
(5.11)

dz is the random DEM error between neighbor cells.

Flux is calculated by the modified DEMOM-Downslope algorithm which starts to calculate the influence matrix for every cell and then sums up the matrixes. For a single influence matrix of cell x_i , the discharge flux of cell x_i ($F_{i,0}$) into the influence matrix is the accumulation rate (\mathring{A}) of the cell multiplied by the surface area of the cell (S). The flux of its downcell ($F_{i,k}$) in the influence matrix accumulates input fluxes from its upcells.

$$j = 0$$
 $F_{i,0} = \overset{\circ}{A}S$ (5.12)

While ice has downcell, that is, ice does not reach the edge of the Antarctica or ice does not flow into a sink, if $x_{i,k}$ is the downcell, the flux for cell $x_{i,k}$ in the influence matrix of x_i ($F_{i,k}$) is given by:

$$j = k \neq 0$$
 $F_{i,k} = \lambda_1 F_{i,k-1}$ $\lambda_1 = \frac{\tan \theta}{2}$ (5.13)

 λ_1 is the partitioning rate from Costa-Cabral and Burges partitioning scheme. Here we just consider the partition of one cardinal direction. For the partition of the other direction, the partitioning rate λ_2 equals:

$$\lambda_2 = 1 - \frac{\tan \theta}{2} \tag{5.14}$$

Once influence matrixes for all cells on the Antarctica are completed, they are added up to get total flux for every cell. The total flux for cell x_j (F_j) is:

$$F_{j} = F_{0,j} + F_{1,j} + \dots + F_{n,j}$$
(5.15)

where $F_{0,j}$, $F_{1,j}$,..., $F_{n,j}$ are the fluxes at cell x_j in all influence matrixes.

So the standard deviations of fluxes are:

$$j = 0$$
 $dF_{i,0} = S \cdot d\overset{0}{A}$ (5.16)

While ice has downcell, that is, ice does not reach the edge of the Antarctica or ice does not flow into a sink, if $x_{i,k}$ is the downcell, the flux error for cell $x_{i,k}$ in the influence matrix of cell x_i ($dF_{i,k}$) is given by:

$$j = k \neq 0$$
 $d^2 F_{i,k} = (\lambda_1 dF_{i,k-1})^2 + (F_{i,k-1}(\frac{d\theta}{\cos^2 \theta}))^2$ $\lambda_1 = \frac{tg\theta}{2}$ (5.17)

Once influence matrixes for all cells on the Antarctica are completed, the total flux error for cell x_j (dF_j) is given by:

$$d^{2}F_{j} = d^{2}F_{0,j} + d^{2}F_{1,j} + \dots + d^{2}F_{n,j}$$
(5.18)

where $dF_{0,j}$, $dF_{1,j}$,..., $dF_{n,j}$ are the flux errors at cell x_j in all influence matrixes.

The balance velocity (V) is:

$$V = \frac{F}{H \cdot (\sin \theta + \cos \theta) \cdot cellsize}$$
(5.19)

So the standard deviation of balance velocity (dV) is:

$$d^{2}V = \left(\frac{dF}{H \cdot (\sin\theta + \cos\theta) \cdot cellsize}\right)^{2} + \left(\frac{FdH}{(\sin\theta + \cos\theta) \cdot cellsize \cdot H^{2}}\right)^{2} + \left(\frac{F \cdot (\cos\theta - \sin\theta)d\theta}{H \cdot cellsize \cdot (\sin\theta + \cos\theta)^{2}}\right)^{2}$$
(5.20)

Figure 5.6 is the random error of balance velocity using influence matrix algorithm. The random error is small. It is about 1% of balance velocity. We apply this error analysis to the test study area discussed in section 5.5.1.



Figure 5.6 Random error of balance velocity using influence matrix algorithm

5.4.2 Errors in algorithmic calculation

It is hard to quantify the errors resulting from algorithms, such as the errors caused by the diffusion problem, cardinal partitioning scheme and cell-size problem. But we can at least estimate where the problems may be more or less severe. For the lost ice problem of cardinal partitioning scheme (also refer to section 4.3.3), we create a map to visualize where the problem happens (Figure 5.7(b)).

Figure 5.7(a) is an influence matrix of a single cell A. The cells with the same colors are processed at the same time – we call the ensemble of cells the front line. Front lines are not determined manually. They arise automatically and somewhat artificially from the process order. Once a front-line cell is processed, we do not allow ice to flow backwards into the front line. In this example, the cell B flow direction points Southwest. According to the cardinal partition scheme, its flux should be partitioned to cell C and D, but cell C is a processed front line. So the part of flux from B which should flow into cell C is routed to cell D. This procedure to adjust flux flow distorts the influence matrix to some degree, but it maintains the total flux. All flow direction arrows in red in Figure 5.7(a) represent adjusted cells. For cell E, its flow direction is in the Northwest, so its flux should be partitioned to cells F and G, but both cells are in a processed front line. Thus the amount of flux from cell E is lost. It is called lost ice. Though we tried to flow the flux from cell E to its neighbor cells which are in the same process line as cell E, in complex terrain, unrealistic flux distribution was created. The black arrows in Figure 5.7(a) indicate the cells that lose ice. Black dots in figure 5.7(b) are the lost ice cells when the modified DEMON-Downslope approach is applied on Antarctica.

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Figure 5.7(a) Conflict of cardinal partition scheme and front line



Figure 5.7(b) Lost ice cells in modified DEMON-Downslope algorithm

5.5 Result validation

We compare the balance velocity result with independently measured velocity to validate our calculation.

5.5.1 Lambert Glacier

Figure 5.8 is the Lambert Glacier mosaic (gray image) showing the location of 1990/1995 GPS traverse data (red dots) collected as part of the Australian National Antarctic Research Expedition.



Figure 5.8 1990/1995 GPS traverse Data



Figure 5.9 Comparison between balance velocities and GPS velocities

Figure 5.9 shows the GPS velocity profiles, balance velocity implementing the Budd and Warner algorithm (DEM Sorting) by ourselves (we do not get their codes or their balance velocities) and balance velocity from the modified DEMON-Downslope algorithm (Influence matrix) along the traverse line. For the data points from 1 to 10, 15 to 20, 65 to 70, these three estimates are similar which validates the balance velocity calculation. For data point 13, the balance velocities do not detect the measured velocity peak, which may simply mean that the feature is too small to be detectable in the 20km cell-size balance velocities. There are three balance velocity peaks where balance velocities are obviously higher than GPS measurements, such as data points from 20 to 25, data point 35 and from 40 to 45. Examining the individual contributing areas for these points, we conclude they are an overestimation by the model rather than disequilibrium of the ice streams. Figure 5.10 is the flow line map (Liu et al. 1999) derived from the DEM along with the locations of GPS measurement points. It is clear that those overestimated velocities are located where flow lines are converging. For those relatively better estimated velocity points, they are located where flow lines are parallel. We conclude that the problem occurs where the flow lines are converging. Consequently the contributing area is likely to be enlarged due to artificial mass diffusion outside of the real flow lines. This is illustrated in Figure 5.11(a). The lines l_1 and l_2 are assumed to be real flow lines. The fluxes of cells B and C flow through cell A because of diffusion. For the parallel flow lines (Figure 5.11(b)), the contributing area of the cell A is less influenced by the diffusion of cells B and C which are outside the real flow lines. We draw the profile on the statistical error (section 5.4). The estimated errors are about 1% of balance velocity

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(Figure 5.12), so current input data for balance velocity estimation are accurate enough.





Figure 5.10 Flow lines (red line) and GPS measurement points (blue dot)





Figure 5.11(b) Contributing area



Figure 5.12 Random error of balance velocity profile

5.5.2 East Antarctic ice stream interferometry velocity

East Antarctic ice stream interferometric velocities (Zhao and Jezek, 2001) are based on the SAR interferometric technique, which is widely used to provide high resolution surface motion vectors. The uncertainty of the flow speed is better than 15 ma^{-1} .

Figure 5.13 is the East Antarctica mosaic (gray image) overlaid with flow stripes coverage (red line) and profile lines (blue line) used in the comparison of the balance and interferometric velocities



Figure 5.13 East Antarctica mosaic overlaid flow stripes and profile lines



Figure 5.14 (a) Profile $A_1 - A_2$

Figure 5.14(a) is the profile line $A_1 - A_2$ for three velocity estimates:

interferometry velocity (blue line), balance velocity using DEM Sorting algorithm (green dash line) and balance velocity using modified DEMOM-Downslope algorithm (red dash line). The interferometry velocity grid cell-size is 500m while for the balance velocities, the cell-size is 20km. For the left part of the profile line (0 km to 200 km) the balance velocities profiles are smoother than that of the interferometry velocities because the small scale tributaries are not detected in the balance velocity model. Past distances of about 370 km, there is a stream which is confined by mountains on two sides. The balance velocities are higher than the interferometry velocity from the distances of about

320 to about 370km because the flow directions at the boundary of the stream make ice flow out of the boundary. There is an obvious offset between influence matrix balance velocities and interferometry velocity at distances of about 370km. This offset is due to the low resolution of DEM which causes the boundary between the stream and a nearby tributary to blur and create an unrealistically wider stream. The DEM Sorting balance velocities are better at distances of about 370km because ice is constraint not to flow higher terrain by DEM Sorting algorithm. The elevations at distances from 300km to 370km are higher than those at distances from 370km to 500km so the ice does not flow out of ice stream in DEM Sorting algorithm. Considering the sum of fluxes across the gate from distances from about 320km to 500km, influence matrix algorithm sum is more similar to interferometry velocity than is the DEM Sorting algorithm sum. Also the balance velocity from DEM sorting algorithm has an unreliable peak at distances 480km, which might be because the DEM sorting algorithm is more affected by DEM errors than the influence matrix algorithm.



Figure 5.14 (b) Profile $B_1 - B_2$

Figure 5.14(b) is the profile $B_1 - B_2$ for the above three velocity estimates. From 0km to 260km, the balance velocities smooth out the small scale tributaries detected by

interferometry velocity due the 20km cell-size, but the velocity magnitudes are similar between balance velocities and interferometry velocity. There is an unreasonable peak at about 160 km in the model calculated with the DEM sorting algorithm. There an unrealistic broadening in balance velocity stream from about 260km to 480km as compared to the interferometry velocity detected stream from 320km to 420km. This is due to the low resolution DEM from distances about 320km and more. Considering the sum of fluxes across the gate from distances about 260km to 480km, the sum of balance velocities is obviously higher than that of interferometry velocity. Mass balance calculations by Zhao (2001) suggest this area may indeed be slightly out of balance.



5.5.3 Flow pattern comparison

Figure 5.15 Flow stripes (red line) and balance velocity map (gray image)
Here the approach is applied to the entire Antarctica. Figure 5.15 is the balance velocity map from modified DEMON-Downslope algorithm (gray-scale image) overlaid with RAMP flow stripes (red lines). The whiter color in the gray image indicates higher velocity; the darker color indicates lower velocity. The flow pattern from the balance velocity map matches the flow stripes. Though part of the flow stripes are used to force the ice to flow according to the flow stripes (West Antarctic Ice Streams, ice streams draining into the Filchner Ice Shelf, Lambert Glacier), the flow pattern of the rest flow stripes independently matches that from balance velocity. Moreover, balance velocity extends the ability to look at the flows from ridges. In the lower right area, there are several obvious ice streams shown by balance velocity but they do not clearly exist in Radarsat Image. At the very least, this observation deserves scrutiny with interferometric velocities to see whether the ice sheet is out of balance or whether the flow stripes are not as reliable indicators of fast glacier flow as thought.

CHAPTER 6

CONCLUSION

We have developed a modified flux estimation approach which is robust against errors in the DEM. Comparisons have been made to existing approaches. The new approach flexibly incorporates refined flow directions in low slope areas which are always problematic for traditional flux calculation algorithms because of small DEM errors. It does not use DEM sorting order to process cells, as does the Budd and Warner algorithm (1996), which makes our algorithm less sensitive to errors in absolution value of DEM measures. Our algorithm is easy to implement and allows us to visualize upslope and downslope areas for an individual cell, which helps to validate results.

We have presented a new technique for incorporating vector information derived from image data into the flux calculation. This new technique is an easy and straight forward method which mitigates flow direction problems from DEM errors in flat or near flat areas.

A new balance velocity model has been structured which compares favorably with other independently derived surface velocities. Differences between balance velocity and measured velocities may be caused by the diffusion problem. Our study indicates where diffusion is more or less of a problem.

Comparing Antarctic Mapping Mission (AMM) flow stripes over the Antarctica with our balance velocity model, in Wilkes Land, we observed agreement in directions between flow stripes and velocities. However the balance velocity model predicts fast glacier flow deeper into interior Antarctica. This is an interesting glaciological observation which will be tested using surface velocity data to be derived from the Modified Antarctic Mapping Mission interferometry campaign (Jezek, in press).

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